

Coupling the dynamics and the molecular chemistry in the Galactic center

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ABSTRACT

Context. Most of the Galactic center (GC) gas moves in nearly circular orbits in a nuclear ring (hereafter Galactic center ring, GCR). This is the case of cloud complexes such as Sgr A or Sgr B, where the gas is dense, warm and exhibits a rich molecular chemistry. The origin of these properties is thought to be some kind of shock, in particular due to the large scale dynamics of the Galaxy. In addition, there are gas clouds moving in highly non-circular orbits known by observations of low density tracers as CO(1-0). The physical conditions of the clouds moving with non-circular velocities are not well-known.

Aims. We have studied the physical conditions of the gas in non-circular orbits with the aim of better understanding the origin of the outstanding physical conditions of the GC molecular gas and the possible effect of the large scale dynamics on these physical conditions.

Methods. Using published CO(1-0) data, we have selected a set of clouds belonging to all the kinematical components seen in the longitude-velocity diagram of the GC. We have done a survey of dense gas in all the components using the $J = 2 - 1$ lines of CS and SiO as tracers of high density gas and shock chemistry.

Results. We have detected CS and SiO emission in all the kinematical components. The gas density and the SiO abundance of the clouds in non-circular orbits are similar those in the GCR. Therefore, in all the kinematical components there are dense clouds that can withstand the tidal shear. However, there is no evidence of star formation outside the GCR. The high relative velocity and shear expected in the dust-lanes along the bar major axis could inhibit the star formation process, as observed in other galaxies. The high SiO abundances derived in the non-circular velocity clouds are likely due to the large-scale shocks that created the dust lanes.

Key words. ISM: kinematics and dynamics – ISM: molecules – Galaxy: center – Galaxies: ISM – Radio lines : ISM

1. Introduction

The longitude-latitude maps of the Galactic center (GC) show several major cloud complexes: the $l = 1.3^\circ$ -complex (a huge molecular complex located at longitudes between $l = 1.2^\circ$ and $l = 1.6^\circ$), the Sgr D complex at $l = 1^\circ$, the Sgr B complex near $l = 0.7^\circ$, the Sgr A cloud near $l = 0^\circ$ and the Sgr C cloud near $l = -0.5^\circ$ (see for instance, Scoville 1972; Liszt & Burton 1978; Bania 1977; Bally et al. 1987; Dahmen et al. 1997; Oka et al. 1998a,b; Bitran et al. 1997). In the longitude-velocity diagram (hereafter *lv*-diagram, see Fig. 1), the Sgr A ... D

complexes are approximately located along a line passing through the origin as expected for gas moving along an almost circular ring. Therefore, in a face-on view of the Galactic center these complexes would be located in a circular ring (Fig. 2). Sofue (1995) proposed that this ring is composed of two arms (Arm I and Arm II). Hereafter, we refer to this structure as the Galactic center ring (GCR). However, the GCR clouds and the $l = 1.3^\circ$ -complex are not the only cloud components in the GC. The *lv*-diagram also exhibits many structures corresponding to clouds moving in non-circular trajectories like the Clump 2, the Connecting Arm, or the structures that we have labeled as J, K, L, M, N, O and P (Fig. 1). The

easiest explanation for the non-circular trajectories is the gas response to a bar potential (Binney et al. 1991; Fux 1999).

In the inner region of a barred galaxy there are two main families of closed orbits in the reference frame of the bar (Contopoulos & Papayannopoulos 1980): the x_1 orbits, which are elongated along the bar major axis and the x_2 orbits, which are elongated along the minor axis of the bar, when there exists inner Lindblad resonances (ILR). Numerical simulations of the gas dynamics (Athanioula 1992; Fux 1999) show that the gas in the center tends to follow nearly circular x_2 orbits creating a nuclear ring (the GCR). At larger radii, beyond the ILRs, the gas tends to follow x_1 orbits along the bar. In fact, the gas clouds are subject to collisions, and do not follow strictly these orbits, but are frequently kicked off into neighboring orbits, in general with lower energy. In particular, at the transition region between the x_1 and x_2 flows, collisions are favored, and there might be a scatter, or “spray” of gas clouds moving along the x_1 flow. The sprayed gas will collide with the material on the opposite side of the bar major axis giving rise to shocks in the molecular cloud component and to the characteristic dust lanes that are observed near the leading edges of the bars in external galaxies. Losing energy and angular momentum, the gas of the dust lanes will progressively spiral down to the nuclear ring and it will fall in after only a few rotations (see Fig. 20 of Fux 1999).

In the GC, Fux (1999) has identified the Connecting Arm as the near-side dust lane. The structure K, which seems to be linked to the Connecting Arm at $l \sim 3^\circ$ (Fig. 1), represents gas clouds that move in the decelerating section of elongated orbits, i.e, gas that does not fall in the nuclear ring at the first passage around the $x_1 - x_2$ interaction region (see Fig. 2). The structure J has nearly the same inclination as that of the Connecting Arm (Fig. 1), and it could be a kind of small dust lane, but in contrast with the Connecting Arm and the structure K, the gas in the structure J will easily interact with the material of the nuclear disk in the area of the $l = 1.3^\circ$ -complex (see Fig. 2). Thus, the $l = 1.3^\circ$ -complex would be the contact point of the dust lane and the GCR as proposed by Hüttemeister et al. (1998) and Fux (1999). The clump at $l = 5^\circ$ and the Clump 2 are gas clouds that are about to enter the dust lane shock. The lv -diagram at $l < 0$ is more difficult to interpret. One should note that the lv -diagram is not expected to be symmetric. Due to perspective effects the far side dust lane is elongated along the line of sight and thus it is a vertical structure in the lv -diagram close to the extremity of the nuclear ring (see Fig. 23 of Fux 1999). The complexes L, N, and P could indeed belong to the far-side dust lane shock. The structure M would be part of the gas spray due to the interaction of the far-side dust lane and the GCR.

The physical conditions of the neutral gas in the GC have been studied by observing high density tracers like CS (Bally et al. 1988; Tsuboi et al. 1999) or HCN (Jackson et al. 1996) and by multi-

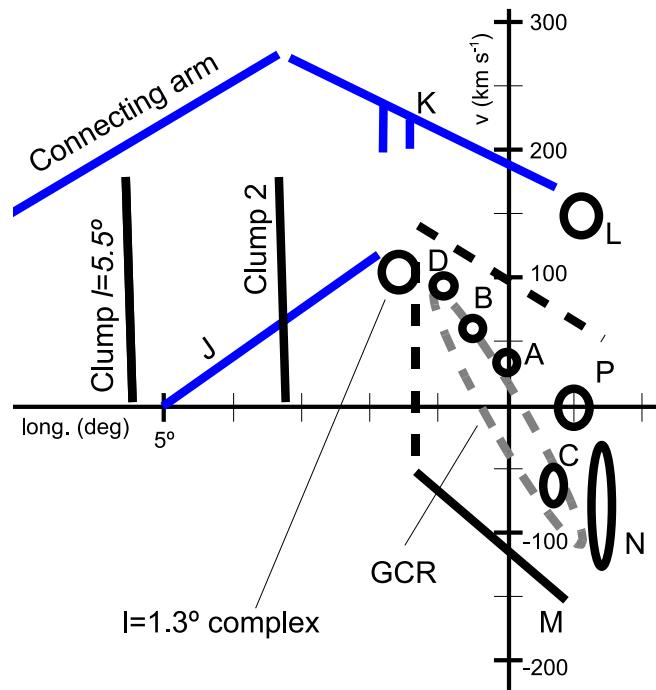


Fig. 1. Schema illustrating the main kinematical structures in the longitude-velocity diagram of the inner degrees of the Galaxy. The labels A to D stand for Sgr A to Sgr D complexes. See the text for an explanation of the other labels

level studies of NH_3 (Hüttemeister et al. 1998), H_2 (Rodríguez-Fernández et al. 2001, 2004) or CO (Sawada et al. 2001; Martin et al. 2004). These works, among others, have shown that the clouds in the GCR and the $l = 1.3^\circ$ -complex are dense (10^4 cm^{-3}) and warm (up to $\sim 150 \text{ K}$). In addition, the GCR clouds present widespread emission and high abundances of molecules like SiO, SO, CH_3OH , or $\text{C}_2\text{H}_5\text{OH}$ (Martín-Pintado et al. 1997, 2001; Lis et al. 2001; Hüttemeister et al. 1998), which are considered good shock tracers. Indeed, shocks are commonly invoked to explain the molecular chemistry and the high temperatures of the neutral gas in the GCR, the $l = 1.3^\circ$ -complex and the Clump 2 (Martín-Pintado et al. 1997, 2001; Hüttemeister et al. 1993, 1998; Lis et al. 2001; Rodríguez-Fernández et al. 2004). The origin of the shocks could be related to Wolf-Rayet winds (Martín-Pintado et al. 1999) or to cloud-cloud collisions induced by the large scale dynamics of the Galaxy (Wilson et al. 1982; Hüttemeister et al. 1998; Martín-Pintado et al. 2001; Sato et al. 2000; Rodríguez-Fernández et al. 2000, 2004).

In contrast to the GCR and the $l = 1.3^\circ$ -complex, little is known about the physical conditions of the GC clouds with non-circular velocities. Up to now, the only large scale maps that have been used to study the physical conditions of the ensemble of the GC molecular gas are those of Oka et al. (1998a) and Sawada et al. (2001). They have mapped the CO(2-1) and CO(1-0) lines and

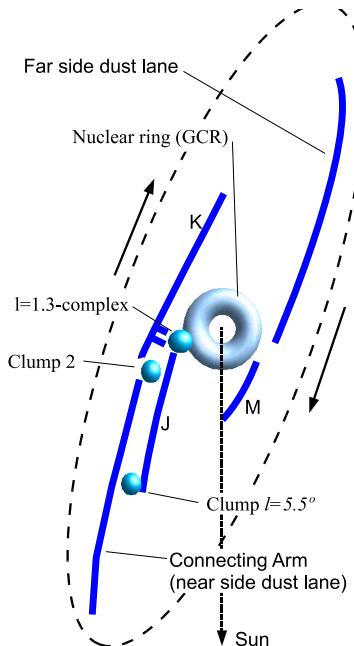


Fig. 2. Schema of a plausible face on view of the Galactic center

found a rather constant line ratio (~ 1) in the GC clouds (including Clump 2). A few clouds of the upper-right corner of the lv -diagram have been observed in SiO and NH₃ by Hüttemeister et al. (1993, 1998) and in H₂ and fine structure lines by Rodríguez-Fernández et al. (2004). Their densities, temperatures and excitation mechanisms seem to be similar to the GCR clouds.

Thus, in contrast to the GCR cloud complexes, the physical conditions of the gas with non-circular velocities that constitute the characteristic shape of the lv -diagram are not well-known. However, the study of the properties of all these kinematical components can give us important hints on the origin of the outstanding physical conditions of the GC clouds and on the possible links between these physical conditions and the global gas dynamics. This is the purpose of this paper. We have selected a sample of clouds belonging to all the kinematic components seen in the GC lv -diagram and we have undertaken a survey of dense gas using CS and SiO as tracers. Due to their high dipole moments, both CS and SiO are high density tracers ($\geq 10^4 \text{ cm}^{-3}$). In addition, the SiO molecule traces dense gas affected by shocks (Martin-Pintado et al. 1992). Since the CS and SiO molecules are quite similar and they have similar energy level distributions, the SiO/CS line ratio is expected to depend mainly on the relative abundance of the two molecules.

2 Observations

We have selected 61 positions in the longitude range $3.5^\circ > l > -1.5^\circ$ using the Bally et al. (1987) CO(1-0) data cube. In fact, due to the crowded velocity fields in the GC, with these 61 pointings we have been able to

study 161 clouds (both GCR and non-GCR clouds). Table 2 presents the Galactic coordinates and velocities of the different clouds. For an easy presentation and discussion of the results we have grouped the clouds that occupy a same region in the (l, b, v) space or that belong to one of the coherent kinematical structures discussed above.

We have observed 43 of the selected positions in the J=2-1 line of CS and SiO using the University of Arizona 12 meter antenna at Kitt Peak in February 2003. The half power beam width of the telescope is 68" and 77" at the frequencies of the CS and SiO lines, respectively. The temperatures, in the T_R^* scale, have been converted to main-beam temperatures (T_{mb}) using the expression $T_{mb} = T_R^*/\eta_m^*$ with corrected main beam efficiencies η_m^* of 0.95 and 0.94 for the SiO and CS spectra, respectively. As backends, we used filter-banks with 1 MHz channels, which provide a velocity resolution of 3.06 and 3.45 km s⁻¹ for the CS and SiO lines, respectively. The *rms* noise in 1 MHz channels is in the range 0.04 – 0.11 K for the CS spectra and in the range 28 – 51 mK for the SiO spectra.

The other 18 positions were observed with the IRAM 30 meter telescope in July 2003. The two 100 GHz receivers were used to simultaneously observe the two lines. The 30m telescope beam size at the frequencies of the CS and SiO lines is 25.5" and 29" respectively. The temperatures, in the T_A^* scale, have been converted to main beam temperatures using the expression $T_{mb} = (F_{eff}/B_{eff}) T_A^*$ with a forward efficiency (F_{eff}) of 0.95 and beam efficiencies (B_{eff}) of 0.78 and 0.77 for the SiO and CS spectra, respectively. As backends, we also used filter-banks with 1 MHz channels. The *rms* noise in 1 MHz channels vary from source to source in the range 15 – 75 mK for the CS spectra and in the range 11 – 57 mK for the SiO spectra. The lowest *rms* are needed to obtain good detections of some clouds with non-circular velocities as those in the structure K. For comparison, the large scale CS(2-1) map of Bally et al. (1987) had a *rms* noise of 0.15 K in 1MHz filters.

We have also observed the ^{13}CO (2-1) and (1-0) lines towards 12 of the positions observed with the 30m telescope (see Table 1). The IRAM 30m beam size at the frequency of these lines is $11''$ and $22''$, respectively. The temperatures have been converted to T_{mb} using forward and beam efficiencies of 0.95 and 0.75 for the ^{13}CO (1-0) line and 0.91 and 0.55 for the ^{13}CO (2-1) line. All the observations were done in position switching mode. The emission free OFF positions have been carefully selected from the Oka et al. (1996) maps and our own spectra of different observational campaigns.

3. Results

Figure 3 shows the CS and SiO spectra obtained towards all the observed positions and Table 2 gives the velocity, widths and intensities of the CS lines as derived from Gaussian fits to the 161 lines detected. When the SiO lines are detected, the velocities and the linewidths are in agree-

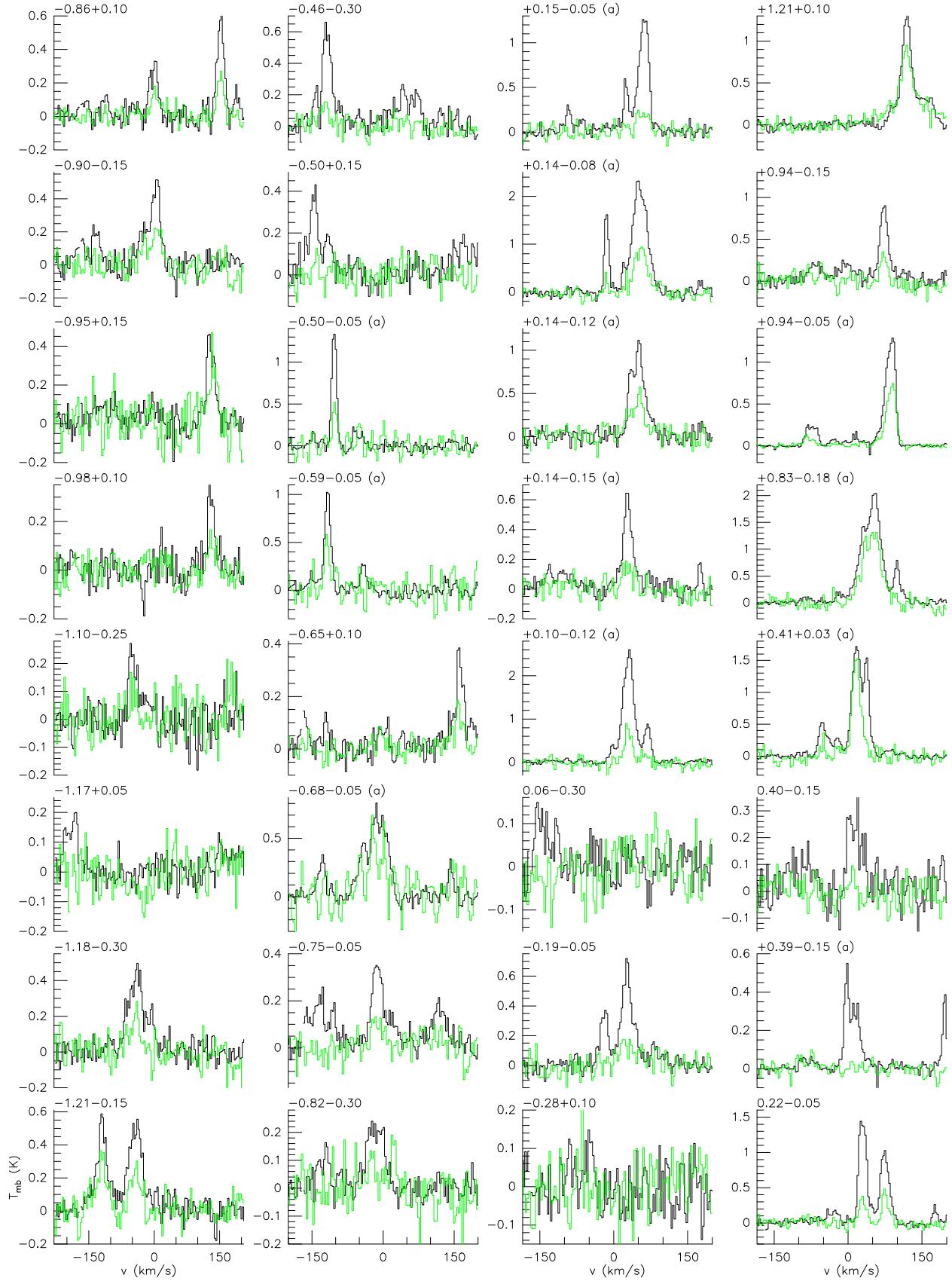


Fig. 3. CS(2-1) spectra (black lines) and SiO(2-1) spectra (multiplied by 2, green lines) towards all the observed positions, which are written in the upper-left corner (in Galactic coordinates). The positions observed with the 30m telescope are indicated with a label (a). The velocity scale is the same for all the spectra in the same column.

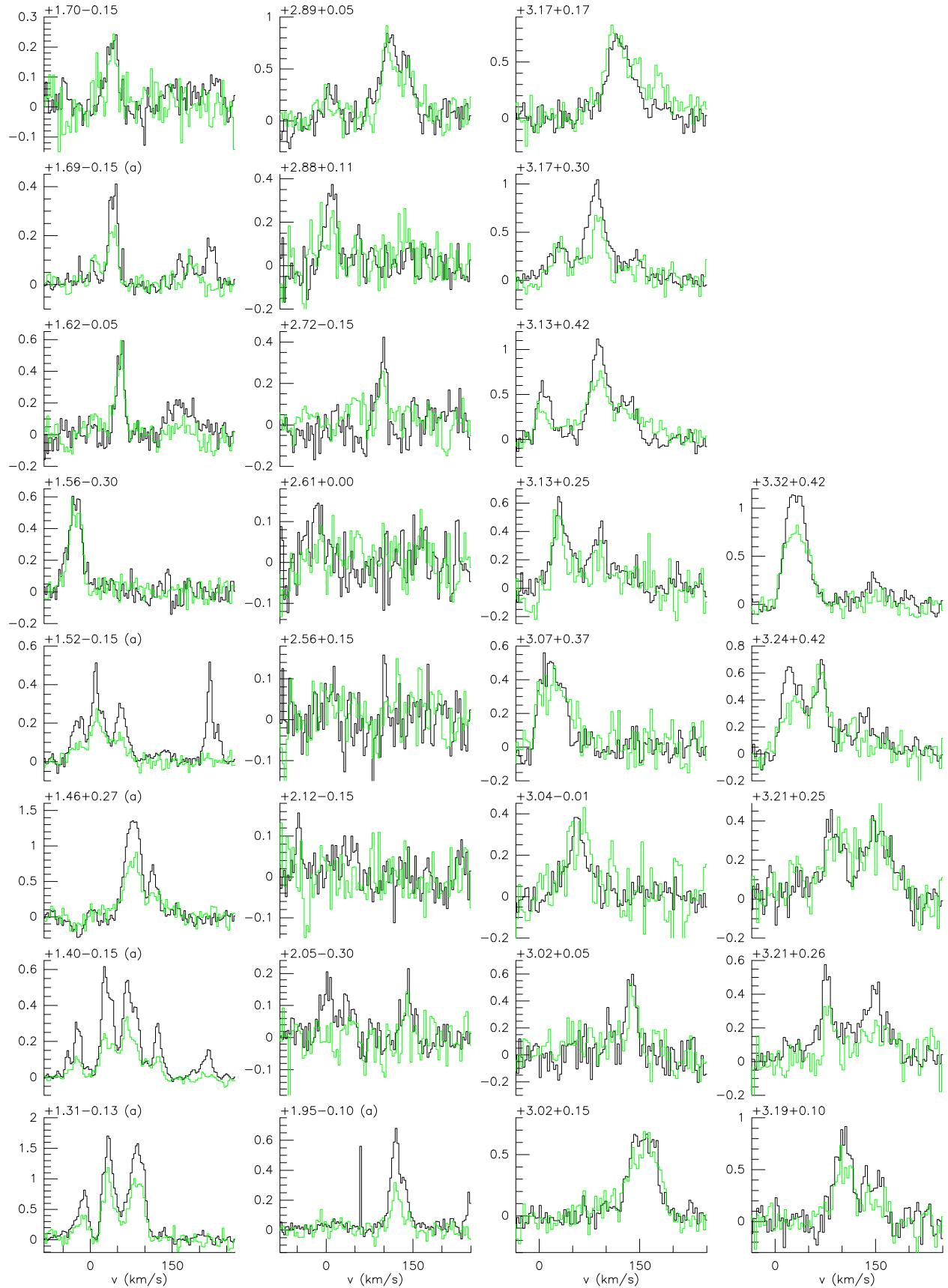
**Fig. 3.** Continuation.

Fig. 4. In grey-levels and contours we show the longitude-velocity diagram obtained with the CO(1-0) data by Bally et al. (1987) at different latitude ranges. We also indicate the different kinematical features discussed in this paper (note that not all the features that are seen in different panels are labeled in all of them). Blue triangles and bars represent the peak velocity and the full width at half maximum (FWHM) of the sources detected in CS. Black points and bars show the peak velocity and FWHM of the sources that have been detected in CS and SiO.

ment within errors with those of the CS lines and Table 2 shows the SiO to CS intensity ratio without any correction by the slightly different telescope beams. For the clouds where the SiO line has not been detected, Table 2 gives upper limits to the SiO(2-1)/CS(2-1) line ratio as derived from the 3σ limits to the SiO intensity. The results are also displayed graphically in Fig. 4. Blue triangles and bars represent the peak velocity and the full width at half maximum (FWHM) of the sources detected in CS. Black points and bars show the peak velocity and FWHM of the sources detected in CS and SiO.

We have detected CS emission in all the kinematical structures of the GC *lv*-diagram (Fig 4, Table 2). The large linewidths measured ($8\text{--}64\text{ km s}^{-1}$) are typical of Galactic center clouds. The CS intensities range between the high values (0.5–1.5 K) measured for the sources in the $l = 1.3^\circ$ -complex, the Clump 2 and the GCR clouds to 0.1–0.3 K in the faint structures that constitute the upper and lower edges of the *lv*-diagram (K, M, R), including the vertical appendix of the component K at $l \sim 1.7^\circ$, which can be the signature of mass transfer from the elongated orbits to the GCR (Fux 1999). The clouds in the right edge of the *lv*-diagram (L, N, P) and those in the structure J exhibit intermediate CS intensities (0.2–0.6 K). The observations presented here demonstrate that there are dense CS-emitting clouds in all the kinematical components seen in the CO *lv*-diagram and not only in the GCR and the $l = 1.3^\circ$ -complex clouds detected by Bally et al. (1987).

We have also detected SiO(2-1) emission in at least a few clouds of all the kinematical components. The SiO(2-1) line has been detected in all the observed clouds of the $l = 1.3^\circ$ -complex and Clump 2. The SiO/CS line ratio is in the range 0.3–0.6. The SiO line has also been detected in two clouds of the component K: one cloud near the CO maximum at $l \sim 1.4$ and one cloud in the vertical appendix at $l \sim 1.7^\circ$. The SiO/CS ratio is 0.14 and 0.5, respectively. SiO emission has also been detected in the lower edge (component M) and the negative longitude extremity of the *lv*-diagram (components P and N, including a few clouds with different velocities along the structure N). The SiO/CS ratios are in the range 0.2–0.5. Similar line ratios have been measured towards several positions in the structures J and L. In most of the clouds where the SiO line has not been detected the CS is indeed rather weak (~ 0.1 K) and the upper limits to the SiO/CS line ratio are not significant.

4. Gas density and cloud stability

To get further insight in the physical conditions in the different kinematical structures we have obtained ^{13}CO (1-0) and (2-1) data for some sources. Table 1 shows ^{13}CO (1-0) line intensities and ^{13}CO (2-1)/(1-0) line ratios without any correction by the different beams. The line ratio is 0.6–1.5 for most of the sources (it would be a factor of two higher if the emission were extended and homogeneous). We have used a Large Velocity Gradient (LVG) code to determine the H_2 density and the ^{13}CO column density for two values of the kinetic temperature (T_K): 20 and 150 K. Table 1 shows the results obtained for a kinetic temperature of 20 K. The measured line ratios imply an H_2 density of $10^{3.5} - 10^{4.5} \text{ cm}^{-3}$ in the GCR clouds and $10^3 - 10^{3.5} \text{ cm}^{-3}$ in the components K or M. This is in perfect agreement with the results of Sawada et al. (2001) and Martin et al. (2004). Assuming a kinetic temperature of 150 K (Hüttemeister et al. 1993; Rodríguez-Fernández et al. 2000, 2001) the H_2 density would be a factor of 3 lower.

In addition, the detection of the CS(2-1) and SiO(2-1) lines in all the velocity components imply the presence of high densities ($\geq 10^4 \text{ cm}^{-3}$, Linke & Goldsmith 1980) in the core of the clouds. Therefore, all the structures seen in the CO *lv*-diagram contain dense molecular gas. In particular, this is true for the gas moving in elongated orbits related to the dust lanes (like the components J and K) or about to enter the dust lanes shocks (Clump 2, see also Stark & Bania 1986). This result is in agreement with what is found in external galaxies, where the density contrast between the dust lanes and the nuclear ring is low (see for instance the results on NGC 5383 by Sheth et al. 2000).

The density of the clouds is directly related to the problem of cloud stability, which is of crucial interest to understand the star formation process. Due to the strong tidal forces in the Galactic center, the clouds must be denser than in the disk of the Galaxy. This question has been studied in number of works (Stark & Bania 1986; Güsten 1989; Stark et al. 1989, 2004). For example, Stark et al. (2004) have found that the minimum density (in cm^{-3}) for a cloud to withstand the tides can be expressed as:

$$n \approx 10^{3.5} \text{ cm}^{-3} \left[\frac{\kappa}{1000 \text{ Gyr}^{-1}} \right]^2 \quad (1)$$

where κ is the epicyclic frequency, which can be estimated to 1036 Gyr^{-1} for a radius of 150 pc and to 506 Gyr^{-1}

Table 1. Summary of the ^{13}CO (2-1) and (1-0) observations. Intensity of the ^{13}CO (1-0) line (T_{mb}) and ^{13}CO (2-1)/(1-0) intensity ratio calculated from the T_{mb} temperatures without any correction for the different beam sizes. When the ^{13}CO (2-1) line has not been detected we have used 3σ upper limits to calculate the 2-1/1-0 ratio. The LVG results for $T_K=20$ K are also displayed. For the sources where the ^{13}CO (2-1)/(1-0) ratio could not be determined, we have assumed a H_2 density of $n = 10^{3-4.5} \text{ cm}^{-3}$. The SiO column density, $N(\text{SiO})$, has been calculated for a H_2 density of $10^{4.5} \text{ cm}^{-3}$. The SiO abundance relative to H_2 , $X(\text{SiO})$, has been calculated using a ^{13}CO abundance of $5 \cdot 10^{-6}$ (see text).

l deg	b deg	v km/s	Feature	^{13}CO (1-0) T_{mb} (K)	2-1/1-0	$\log(n(\text{H}_2))$ $\log(\text{cm}^{-3})$	$N(^{13}\text{CO})$ 10^{16} cm^{-2}	$N(\text{SiO})$ 10^{13} cm^{-2}	$X(\text{SiO})$ 10^{-9}
-0.5030	-0.0500	144.5	L	0.8±0.2	0.7±0.4	3.0 – 4.0	1.1	< 1.1	< 5
-0.5900	-0.0500	152.8	L	0.5±0.2	< 1.7	< 4.3	0.38	< 0.8	< 9.9
-0.6880	-0.0500	142.5	L	1.0±0.7	0.7±0.5	3. – 3.5	0.6	1.0	7.9
0.9452	-0.0500	86.2	L	4.5±0.3	1.0±0.1	3.4 – 3.6	12	7.3	3.0
0.1500	-0.0500	154.7	K	1.2±0.1	0.6±0.1	3 – 3.5	0.45	< 0.3	< 3.1
1.3100	-0.1300	204.3	K	1.1±0.3	2 – 4	< 0.2	< 0.33
0.3899	-0.1500	195.8	K	0.8±0.3	0.8	< 2.5	< 15
1.4027	-0.1500	214.3	K	1.0±0.1	0.6±0.1	3. – 3.4	2	0.4	0.88
1.5220	-0.1500	219.5	K	1.6±0.3	2.5	0.2	0.4
0.3899	-0.1500	-92.2	M	0.5±0.3	1.3±0.8	3.5 – 4.3	0.2 – 1	0.3	1.3 – 6.8
0.3899	-0.1500	-71.5	M	0.3±0.2	< 2.3	< 4.3	> 0.6	0.6	< 5
0.9452	-0.0500	-81.0	M	1.0±0.3	1.0±0.4	3. – 3.8	1.1	0.7	3.2
0.9452	-0.0500	-65.9	M	1.6±0.3	0.7±0.2	3. – 3.5	2.8	0.4	0.63
-0.6880	-0.0500	-41.7	P	0.9±0.4	1.5±0.7	3.5 – 5	1.7	2.7	7.9
-0.6880	-0.0500	-17.5	P	1.0±0.4	2.2±0.9	3.5 – 5	1.0	3.2	16
-0.6880	-0.0500	2.3	P	2.0±0.4	1.5±0.3	3.5 – 5	1.0	5.0	25
-0.5900	-0.0500	-41.3	P	1.6±0.2	0.9±0.2	3.5 – 4	2.2	< 1.4	< 3.2
1.3100	-0.1300	-13.8	1.3	2.0±0.3	0.8±0.2	3.5 – 4.3	5.7	0.2	0.16
1.3100	-0.1300	33.7	1.3	1.9±0.3	1.0±0.2	3.5 – 4.3	5.7	9.5	8.3
1.3100	-0.1300	85.0	1.3	2.4±0.3	1.1±0.2	3.5 – 4.3	7.5	5.3	3.5
1.3100	-0.1300	135.0	1.3	0.5±0.3	< 1.7	< 10^4	0.4	< 0.1	< 1.3
1.4027	-0.1500	-41.8	1.3	1.1±0.1	0.6±0.1	3.	0.6	< 0.1	< 1.0
1.4027	-0.1500	-22.9	1.3	1.5±0.3	0.7±0.2	3 – 3.5	1.9	0.5	1.2
1.4027	-0.1500	29.7	1.3	1.5±0.2	0.9±0.2	3.3 – 3.5	2.2 – 3.5	1.4	2.5
1.4027	-0.1500	72.3	1.3	1.6±0.2	0.9±0.1	3.3 – 3.5	6.6	2.3	1.7
1.4027	-0.1500	123.7	1.3	0.6±0.0	1.3±0.1	3.5 – 4.0	0.85	0.5	3
1.5220	-0.1500	-23.4	GCR	1.1±0.8	0.8±0.6	3.2 – 4	2.3	0.5	1.0
1.5220	-0.1500	56.0	GCR	1.1±0.8	0.9±0.7	< 4.3	3.4	0.9	1.3
-0.6880	-0.0500	-129.1	GCR	2.6±0.4	1.2±0.2	3.5 – 4.5	4.5	1.1	1.2
-0.5900	-0.0500	132.9	GCR	0.3±0.0	< 3.3	...	16	< 0.8	< 0.25
-0.5900	-0.0500	-148.5	GCR	1.1±0.2	1.0±0.3	3.5 – 4.5	2.5	< 2.6	< 5.2
-0.5900	-0.0500	-116.7	GCR	5.3±0.2	1.0±0.1	> 3.5	9.5	3.0	1.6
-0.5030	-0.0500	-147.4	GCR	1.1±0.2	0.6±0.3	2.8 – 3.5	0.5	0.1 – 0.8	1.3 – 8
-0.5030	-0.0500	-53.5	GCR	1.0±0.2	1.3±0.4	3.5 – 4.5	1.4	0.7 – 2.1	2.5 – 7.5
0.1500	-0.0500	25.5	GCR	4.0±0.4	1.2±0.2	> 3.5	6	< 1	0.79
0.1500	-0.0500	61.1	GCR	5.3±0.4	1.4±0.1	> 3.5	14	1.5	0.54
0.8300	-0.1800	31.3	GCR	5.0±0.5	1.0±0.1	> 3.5	16	6.4	2
0.8300	-0.1800	56.9	GCR	0.5±0.5	1.3±1.5	3.5 – 4.5	1.1	8.5	30
0.8300	-0.1800	99.1	GCR	1.8±0.5	1.0±0.3	3.2 – 4.5	1.4	1.4	5
0.4100	0.0300	-51.0	GCR	1.1±0.5	1.2±0.6	3.5 – 4.2	1.9	2.5	6.6
0.4100	0.0300	-21.9	GCR	1.5±0.5	0.8±0.3	3. – 3.5	3.3	1.1	1.7
0.4100	0.0300	17.0	GCR	3.9±0.5	1.0±0.2	3.5 – 4.3	6.8	12	8.8
0.4100	0.0300	37.9	GCR	5.1±0.5	1.1±0.1	> 3.5	8.2	< 1.3	< 0.79
0.1000	-0.1200	-2.0	GCR	0.8±0.0	2.2±0.7	> 3.5	1.1	0.9	4
0.1000	-0.1200	30.1	GCR	6.0±0.6	1.0±0.2	3.5 – 4.3	25	8.2	1.64
0.1000	-0.1200	68.1	GCR	1.9±0.6	1.4±0.6	> 3.5	3.8	0.9	1.2
0.3899	-0.1500	-3.5	GCR	1.8±0.2	1.5±0.2	> 3.7	1.9	< 0.3	< 0.7

for a radius of 450 pc (Stark et al. 2004). Assuming a sun-

GC distance of 8.5 kpc, 150 pc is the radius of the GCR.

Therefore, the critical density is $\sim 3 \cdot 10^3 \text{ cm}^{-3}$. The clouds with non-circular velocities are located at larger galactocentric radius than the GCR, therefore the critical density is lower. For instance, the critical density for a cloud located at 450 pc is a factor of ~ 4 lower than for a cloud located at 150 pc (Stark et al. 2004). Therefore, the CS and SiO clouds in non-circular orbits are gravitationally bound and the dust lanes do not contain diffuse gas only (see also Stark et al. 2004).

However, there are no signs of star formation in the dust lanes. The surveys of hydrogen recombination lines (at centimeter wavelengths) in the inner 2 degrees (Pauls & Mezger 1975) show that the HII regions are mostly concentrated along a line that runs through the origin of the lv plane indicating that the star formation is confined to the GCR (see Fig. 5). We have also looked for HII regions at larger longitudes in the catalog by Paladini et al. (2003). Figure 5 shows that there are two HII regions at longitude $l \sim -0.6^\circ$ associated with the structure P (or to the spiral arms in the line of sight) and one ionized region at $l = -1.4^\circ$ associated with Sgr E (and therefore to the GCR). At longitudes $l > 1.2^\circ$ all the known HII regions are most likely associated with the spiral arms except one source with $l = 2.6^\circ$ and $v = 102 \text{ km s}^{-1}$ that could be associated with the structure J. However this structure is basically seen at latitudes $b < -0.1^\circ$ (see Fig. 4) and the HII region is located at a latitude $b = 0.1^\circ$ (Paladini et al. 2003). Therefore, the data available show no signs of star formation in the cloud complexes associated to the dust lanes (Clump 2, J, K, M). However, we can not rule out the presence of ultracompact HII regions (UC-HII) embedded in the dense clumps since very young HII regions could be optically thick at centimeter wavelengths. The presence of isolated UC-HII regions is a possibility that can only be ruled out by deep high resolution searches for recombination line emission at millimeter wavelengths. Nevertheless, such very young UC-HII regions are usually associated with more evolved HII regions that should have been detected by Pauls & Mezger (1975) and the surveys compiled by Paladini et al. (2003). In summary, with the available recombination line data there are no signs of massive star formation in the dust lanes.

Indeed, the lack of star formation along the dust lanes is common in the observations of early-type barred galaxies, where there is star formation in the nuclear ring and at the extremities of the bar, but there are no signs of star formation in the dust lanes along the bar (see for instance Tubbs 1982; Reynaud & Downes 1998; Sheth et al. 2000). The inhibition of the star formation can be an effect of the high velocities, relative to the gas in the dust lanes, of the dense clouds whose orbits intersect the dust lanes (Tubbs 1982). These velocities are much higher in barred than in normal galaxies. The models by Tubbs (1982) showed that the clouds that enter the dust lanes at velocities higher than a critical value ranging between 20 and 60 km s^{-1} must be inhibited from forming stars to explain the observations of NGC 5383. The high ve-

locity can cause a quick compression of the clouds followed by a rapid expansion which disperses the majority of the cloud. This mechanism has also been invoked by Reynaud & Downes (1998) to explain the observations of NGC 1530. This mechanism could also inhibit the star formation in the gas with non-circular velocities in the GC.

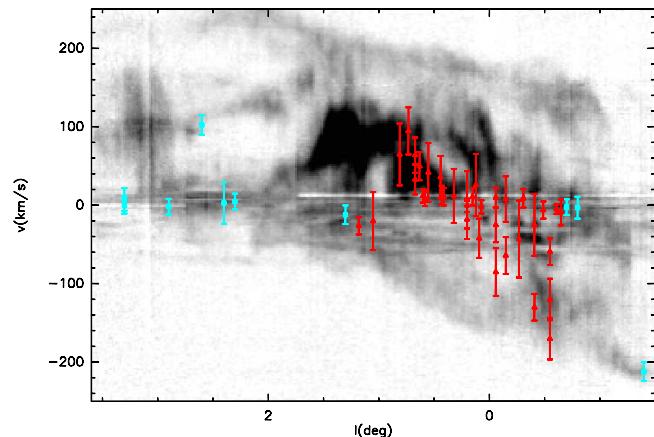


Fig. 5. In grey-levels we show the longitude-velocity diagram obtained integrating the CO(1-0) data by Bally et al. (1987) in the latitude range $-0.3^\circ < b < 0.3^\circ$. Red triangles and bars represent the peak velocity and the full width at half maximum (FWHM) of the radio recombination lines detected in the longitude range $1.18^\circ > l > -0.55^\circ$ (Pauls & Mezger 1975). Blue squares and bars represent the peak velocity and the FWHM of the radio recombination lines detected at $l > 1.18^\circ$ and $l < -0.55^\circ$ (Paladini et al. 2003).

5. SiO abundances and the origin of the SiO emission

As already mentioned, the SiO/CS line ratio is expected to depend mainly on the relative abundance of the two molecules. For example, in the cloud sample studied by Hüttemeister (1993) and Hüttemeister et al. (1998), which is basically composed by clouds of the GCR and the $l = 1.3^\circ$ -complex, the SiO abundance shows a clear trend to increase from $2 \cdot 10^{-10}$ to 10^{-8} with increasing SiO(2-1)/CS(2-1) line ratios in the range 0.1-0.6. These should also be the SiO abundances of the sources of this paper with detected SiO.

In fact, we have derived the SiO abundances for some sources of our sample using the ^{13}CO data and LVG calculations. Table 1 shows the SiO column densities obtained with the typical temperature ($T_K = 20 \text{ K}$) and density ($n(H_2) \sim 10^{4.5} \text{ cm}^{-3}$) of the SiO clouds in the GC (Martín-Pintado et al. 1997; Hüttemeister et al. 1998). The abundance of ^{13}CO relative to H_2 is $5 \cdot 10^{-6}$ in the Galactic center (see Hüttemeister et al. 1998, and references therein). Therefore, the SiO to ^{13}CO column

density ratios imply abundances of SiO relative to H₂ of (0.5–8) 10^{−9} for the sources with detected SiO, including those in the structures K and M. These abundances are in agreement with the results of Hüttemeister et al. (1998) and Martín-Pintado et al. (1997, 2000) for clouds located mainly in the GCR and the $l = 1.3^\circ$ -complex. For one source in the Sgr B2 area ($l = 0.83^\circ$, $b = -0.18^\circ$, $v = 57$ km s^{−1}) and two in the feature P we have even measured a SiO abundance of $\sim 10^{-8}$, which is comparable to the highest abundances derived by Hüttemeister et al. (1998) in the $l = 1.3^\circ$ -complex.

The SiO abundance shows extreme variations in the Galaxy. The abundance is $< 5 \cdot 10^{-12}$ in dark clouds (Ziurys et al. 1989) and $10^{-11} - 10^{-10}$ in diffuse translucent clouds and in spiral arms clouds (Greaves et al. 1996; Turner 1998; Lucas & Liszt 2000). Schilke et al. (2001) have measured an abundance of $\sim 10^{-11}$ in the photo-dissociation region (PDR) of the Orion bar and S140. In contrast, the SiO abundance in shocked regions associated to star formation is $\sim 10^{-8}$ (Martín-Pintado et al. 1992; Bachiller & Perez Gutierrez 1997). In these regions the Si and SiO abundances in the gas phase are increased by the processing of the dust grains by shock waves (Schilke et al. 1997). Thus, the SiO abundances in the non-circular velocity GC clouds measured in this work and those derived in previous works for other GC clouds (Martín-Pintado et al. 1997, 2000; Hüttemeister et al. 1998), are close to the values found in gas shocked by molecular outflows from protostars. Measurements of the SiO abundances in the nuclei of external galaxies like NGC 253, M82, NGC 1068 or IC 342 (García-Burillo et al. 2000, 2001; Usero et al. 2004, 2005) give values from 10^{-10} to a few 10^{-9} , comparable to those measured in the GC.

Two main scenarios have been invoked to explain the widespread high abundances of SiO in the GC and the center of the nearby galaxies cited above:

(i) Grain sputtering by shocks waves of different origins (Martín-Pintado et al. 1997). Hüttemeister et al. (1998) have found evidence of a link between the large scale Galactic dynamics and the origin of the shocks in the GC (high SiO abundances in the $l = 1.3^\circ$ -complex, which they proposed to be due to the interaction of x_1 and x_2 orbits). García-Burillo et al. (2000) and Usero et al. (2005) also use shocks related to the galactic dynamics to explain the SiO emission in NGC 253 and IC 342.

(ii) More recently, a second mechanism has been proposed for regions in the vicinity of X-ray sources: the destruction of small silicate grains by energetic X-rays. This mechanism could play a role in the GCR (Martín-Pintado et al. 2000) and close to the AGN in NGC 1068 (Usero et al. 2004).

As it is well known, the arena of energetic phenomena is in the inner regions of the GC (birth and death of massive stars, energetic radiation, shocks,...) making it very complex to interpret the SiO observations of the GCR clouds. In contrast, as discussed above, there are no signs of star formation in the dust lanes and the SiO abundances are as high as in the other GC sources. The high

SiO abundances measured in the GC clouds with non-circular velocities can be due to the strong shocks that created the dust lanes. This scenario has also been proposed by Meier & Turner (2005) to interpret the detection of HNCO and CH₃OH emission along the bar of IC 342. Usero et al. (2005) have also detected SiO emission in the dust lane along the bar of IC 342 and they have measured high SiO abundances of a few 10^{-9} (about a factor of 5-10 higher than in the nuclear ring). They interpret their findings as the effect of large-scale shocks driven by cloud-cloud collisions along the potential well of the bar.

6. Conclusions

We have shown the presence of dense gas in all the kinematical components of the Galactic center, including the structures with non-circular velocities like the dust lanes along the bar. The gas density of these clouds are higher than the critical density for cloud stability. However, in the GC there are no signs of star formation except in the nuclear ring (the GCR). The high relative velocities and shear expected in other kinematical structures, in particular in those associated with the dust lanes, could inhibit the star formation. This is the mechanism proposed to explain the lack of star formation in the dust lanes along the bar of Galaxies like NGC 5383 and NGC 1530. The high measured abundances of one of the most representative molecules of the GC shock chemistry (SiO) in the cloud components associated with the dust lanes and elongated orbits supports this scenario. Furthermore, it gives us important hints on the origin of the emission of this molecule. The origin of the shocks and therefore the GC molecular chemistry can indeed be strongly influenced by the global dynamics of the Galaxy. More observations of the molecular chemistry in face-on barred galaxies would be very interesting for a better understanding of the links between the gas dynamics and the molecular chemistry.

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Table 2. Galactic longitude and latitude (in degrees), velocity (km s^{-1}), widths (km s^{-1}), and intensities (K) of the CS lines. In addition, the SiO to CS intensity ratio is also shown. The numbers in parentheses are the uncertainties on the last digit. The positions observed with the IRAM 30m telescope are indicated with a superscript. All the other positions have been observed with the 12m telescope.

l	b	v	Δv	I(CS)	$\frac{I(\text{SiO})}{I(\text{CS})}$
l=1.3-complex region					
+1.403 ^a	-0.150	29(1)	22(1)	0.58(3)	0.19(2)
+1.310 ^a	-0.130	33(1)	23(1)	1.58(9)	0.33(3)
+1.403 ^a	-0.150	123(2)	16(1)	0.36(3)	0.17(3)
+1.403 ^a	-0.150	72(1)	33(1)	0.49(3)	0.33(3)
+1.310 ^a	-0.130	85(1)	32(1)	1.53(9)	0.33(3)
+1.310 ^a	-0.130	135(1)	11(4)	0.16(9)	≤ 0.42
+1.218	+0.100	119(4)	56(7)	0.39(3)	0.63(8)
+1.218	+0.100	119(2)	17(1)	0.91(3)	0.24(3)
+1.218	+0.100	163(3)	28(4)	0.26(3)	0.34(11)
+1.460 ^a	+0.270	78(1)	32(5)	1.43(7)	0.28(2)
+1.460 ^a	+0.270	116(2)	16(1)	0.64(7)	0.25(6)
+1.698 ^a	-0.150	3(3)	15(3)	0.09(3)	0.7(3)
+1.522 ^a	-0.150	11(2)	24(1)	0.41(4)	0.27(4)
+1.560	-0.300	-26(1)	29(1)	0.54(4)	0.46(6)
+1.522 ^a	-0.150	-23(2)	23(1)	0.22(4)	0.18(6)
+1.403 ^a	-0.150	-41(1)	8(1)	0.15(3)	≤ 0.11
+1.403 ^a	-0.150	-22(2)	15(1)	0.30(3)	0.20(4)
+1.310 ^a	-0.130	-13(2)	23(1)	0.73(9)	0.25(7)
+1.620	-0.050	152(5)	30(8)	0.10(6)	0.6(6)
+1.708	-0.150	40(2)	25(2)	0.20(2)	0.57(9)
+1.698 ^a	-0.150	42(3)	19(3)	0.39(3)	0.31(4)
+1.522 ^a	-0.150	56(2)	26(1)	0.27(4)	0.22(5)
+1.620	-0.050	55(5)	16(1)	0.54(5)	0.57(7)
Clump 2					
+3.329	+0.425	30(1)	33(1)	1.22(6)	0.34(3)
+3.329	+0.425	150(1)	32(4)	0.24(6)	0.3(2)
+3.235	+0.425	68(3)	22(3)	0.59(6)	0.47(7)
+3.235	+0.425	26(3)	31(3)	0.63(6)	0.32(5)
+3.219	+0.250	84(1)	34(3)	0.38(6)	0.30(10)
+3.219	+0.250	152(1)	58(4)	0.33(6)	0.47(13)
+3.210	+0.263	77(1)	27(2)	0.45(6)	0.32(8)
+3.210	+0.263	146(1)	55(3)	0.33(6)	0.29(11)
+3.193	+0.100	101(1)	27(2)	0.78(6)	0.35(7)
+3.193	+0.100	142(1)	27(3)	0.46(6)	0.22(11)
+3.172	+0.175	119(1)	45(1)	0.72(5)	0.44(7)
+3.172	+0.175	183(2)	31(6)	0.16(5)	0.9(4)
+3.171	+0.300	29(1)	29(2)	0.33(5)	0.41(12)
+3.171	+0.300	84(2)	33(1)	0.92(5)	0.32(4)
+3.171	+0.300	143(1)	37(4)	0.19(5)	0.3(2)
+3.138	+0.425	7(1)	23(2)	0.54(4)	0.32(5)
+3.138	+0.425	88(2)	28(1)	1.03(4)	0.34(3)
+3.138	+0.425	130(1)	30(2)	0.46(4)	0.27(6)
+3.136	+0.250	131(4)	32(7)	0.11(5)	0.4(4)
+3.136	+0.250	88(1)	30(3)	0.35(5)	0.29(12)
+3.136	+0.250	31(2)	29(1)	0.53(5)	0.47(9)
+3.074	+0.375	18(2)	37(1)	0.52(6)	0.38(8)
+3.047	-0.013	59(1)	34(2)	0.35(5)	0.4(2)
+3.020	+0.150	156(1)	46(1)	0.67(6)	0.45(7)
+3.020	+0.050	136(1)	17(1)	0.62(4)	0.36(6)
+2.895	+0.050	116(1)	56(3)	0.74(2)	0.39(2)

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Table 2. Continuation.

<i>l</i>	<i>b</i>	<i>v</i>	Δv	I(CS)	$\frac{I(SiO)}{I(CS)}$
Galactic center ring					
Sgr E					
-1.178	+0.050	-182(2)	21(8)	0.15(2)	≤ 0.87
-1.178	+0.050	-203(2)	14(3)	0.16(2)	≤ 0.87
Sgr B					
+0.830 ^a	-0.180	31(2)	32(3)	1.12(7)	0.29(5)
+0.830 ^a	-0.180	56(1)	27(1)	1.71(7)	0.32(3)
Sgr A					
+0.227	-0.050	74(2)	21(1)	0.97(5)	0.21(2)
+0.150 ^a	-0.050	25(2)	13(1)	0.52(6)	≤ 0.19
+0.150 ^a	-0.050	61(2)	23(2)	1.30(6)	0.08(1)
+0.145 ^a	-0.080	52(1)	36(2)	2.21(9)	0.20(2)
+0.145 ^a	-0.150	28(1)	22(1)	0.59(3)	0.12(4)
+0.145 ^a	-0.120	32(2)	8(1)	0.44(7)	0.30(7)
+0.145 ^a	-0.120	50(1)	29(1)	0.97(7)	0.25(3)
+0.100 ^a	-0.120	30(3)	25(3)	2.53(7)	0.15(1)
+0.100 ^a	-0.120	68(3)	15(3)	0.85(7)	≤ 0.16
Other sources					
+0.830 ^a	-0.180	99(1)	17(1)	0.67(7)	≤ 0.20
+0.945	-0.150	73(2)	17(1)	1.18(6)	0.21(2)
+0.945 ^a	-0.050	86(2)	23(2)	1.31(8)	0.28(2)
+0.410 ^a	+0.030	17(3)	20(3)	1.84(4)	0.41(2)
+0.410 ^a	+0.030	37(3)	13(3)	1.38(4)	≤ 0.10
+0.410 ^a	+0.030	-51(3)	19(3)	0.52(4)	0.35(7)
+0.410 ^a	+0.030	-21(3)	21(3)	0.26(4)	0.3(2)
+0.401	-0.150	12(1)	35(3)	0.24(4)	≤ 0.22
+0.390 ^a	-0.150	-3(1)	9(2)	0.42(2)	≤ 0.11
+0.390 ^a	-0.150	13(1)	28(4)	0.30(2)	≤ 0.15
+0.227	-0.050	29(1)	18(2)	1.51(5)	0.10(1)
+0.227	-0.050	-30(2)	28(4)	0.14(5)	≤ 0.51
+0.150 ^a	-0.050	-92(2)	13(2)	0.27(6)	≤ 0.33
+0.150 ^a	-0.050	-65(1)	14(3)	0.14(6)	≤ 0.54
+0.145 ^a	-0.080	-14(1)	10(1)	1.68(9)	≤ 0.08
+0.100 ^a	-0.120	-2(3)	14(3)	0.35(7)	≤ 0.35
-0.198	-0.050	27(2)	27(1)	0.63(5)	0.11(3)
-0.198	-0.050	85(3)	33(8)	0.10(5)	0.5(3)
-0.198	-0.050	-18(1)	21(2)	0.35(5)	≤ 0.17
-0.460	-0.300	45(4)	34(9)	0.18(4)	≤ 0.25
-0.460	-0.300	72(2)	18(4)	0.18(4)	≤ 0.25
-0.503 ^a	-0.050	-53(1)	21(2)	0.16(4)	0.9(5)
-0.503 ^a	-0.050	-147(1)	8(3)	0.17(4)	0.6 (5)
-0.503 ^a	-0.050	-103(1)	11(2)	1.42(4)	0.18(5)
-0.590 ^a	-0.050	-148(2)	25(6)	0.12(2)	≤ 1.24
-0.590 ^a	-0.050	-116(2)	14(1)	1.02(2)	0.28(1)
-0.688 ^a	-0.050	-129(1)	17(2)	0.31(4)	≤ 0.61
-0.754	-0.050	-132(2)	61(4)	0.17(4)	≤ 0.25
F					
-0.655	+0.100	-165(2)	20(4)	0.14(3)	≤ 0.24
-0.655	+0.100	-117(2)	15(4)	0.10(3)	≤ 0.33
-0.503 ^a	+0.150	-145(6)	21(2)	0.59(5)	≤ 0.17
-0.503 ^a	+0.150	-112(1)	20(3)	0.28(5)	≤ 0.33
-0.285	+0.100	-87(4)	20(9)	0.07(3)	≤ 0.77
-0.285	+0.100	-51(4)	22(7)	0.09(2)	≤ 0.72

Table 2. Continuation.

<i>l</i>	<i>b</i>	<i>v</i>	Δv	I(CS)	$\frac{I(SiO)}{I(CS)}$
L					
-0.503 ^a	-0.050	144(3)	17(4)	0.07(4)	≤ 1.26
-0.590 ^a	-0.050	132(1)	8(1)	0.12(2)	≤ 1.24
-0.590 ^a	-0.050	152(1)	7(2)	0.10(2)	≤ 1.44
-0.688 ^a	-0.050	142(1)	10(1)	0.30(4)	≤ 0.63
-0.754	-0.050	118(2)	41(5)	0.16(4)	≤ 0.26
-0.503 ^a	+0.150	171(4)	52(9)	0.12(4)	≤ 0.63
-0.655	+0.100	161(6)	20(2)	0.29(3)	0.21(5)
-0.862	+0.100	152(1)	18(1)	0.54(4)	0.24(2)
-0.950	+0.150	129(1)	23(2)	0.42(5)	0.40(10)
-0.982	+0.100	130(2)	18(2)	0.30(3)	0.27(5)
O					
+2.895	+0.050	8(2)	24(4)	0.30(2)	≤ 0.38
+2.880	+0.113	9(1)	22(2)	0.32(4)	0.28(1)
+2.720	-0.150	9(4)	19(6)	0.07(4)	≤ 0.55
+2.610	+0.000	-9(2)	28(5)	0.14(3)	≤ 0.47
+2.610	+0.000	-39(3)	16(6)	0.07(3)	≤ 0.78
+2.568	+0.150	4(2)	23(4)	0.12(3)	≤ 0.60
+2.121	-0.150	-47(1)	13(3)	0.14(3)	≤ 0.37
+2.121	-0.150	25(4)	49(8)	0.06(3)	≤ 0.68
+2.056	-0.300	1(2)	23(2)	0.17(3)	≤ 0.31
+2.056	-0.300	36(1)	18(3)	0.11(3)	≤ 0.44
P					
-0.590 ^a	-0.050	-41(1)	14(2)	0.28(2)	≤ 0.59
-0.688 ^a	-0.050	-41(1)	26(3)	0.42(4)	≤ 0.47
-0.688 ^a	-0.050	-17(1)	16(2)	0.57(4)	0.51(7)
-0.688 ^a	-0.050	2(1)	25(2)	0.59(4)	0.47(7)
-0.754	-0.050	-10(2)	37(3)	0.32(4)	0.15(4)
-0.829	-0.300	-18(1)	40(2)	0.22(3)	≤ 0.35
-0.906	-0.150	-2(2)	39(2)	0.40(5)	0.24(4)
-0.655	+0.100	2(2)	24(7)	0.09(3)	≤ 0.35
-0.862	+0.100	-1(1)	23(3)	0.25(4)	0.24(5)
J					
+2.720	-0.150	96(1)	23(3)	0.29(4)	0.45(8)
+2.056	-0.300	141(2)	23(2)	0.16(3)	≤ 0.33
+1.958 ^a	-0.100	122(3)	27(1)	0.56(6)	0.25(3)
K					
+1.958 ^a	-0.100	247(1)	9(3)	0.22(6)	≤ 0.32
+1.708	-0.150	221(4)	31(6)	0.08(2)	≤ 0.85
+1.698 ^a	-0.150	220(3)	20(3)	0.16(3)	≤ 0.20
+1.522 ^a	-0.150	219(3)	16(1)	0.42(4)	≤ 0.09
+1.403 ^a	-0.150	214(6)	21(5)	0.14(3)	0.14(4)
+1.310 ^a	-0.130	204(2)	19(4)	0.11(9)	≤ 0.53
+0.401	-0.150	192(3)	28(8)	0.08(4)	≤ 0.49
+0.390 ^a	-0.150	195(2)	8(1)	0.39(2)	≤ 0.12
+0.145 ^a	-0.150	175(2)	7(1)	0.20(3)	≤ 0.38
+1.698 ^a	-0.150	185(3)	12(3)	0.11(3)	0.5(2)
+1.698 ^a	-0.150	163(3)	13(3)	0.12(3)	≤ 0.25
+1.708	-0.150	157(6)	48(9)	0.53(2)	≤ 0.01
+0.227	-0.050	176(1)	10(1)	0.20(5)	≤ 0.39
+0.150 ^a	-0.050	154(1)	4(1)	0.20(6)	≤ 0.42
+0.145 ^a	-0.080	177(2)	12(4)	0.21(9)	≤ 0.47

Table 2. Continuation.

<i>l</i>	<i>b</i>	<i>v</i>	Δv	I(CS)	$\frac{I(SiO)}{I(CS)}$
M					
+0.401	-0.150	-102(4)	53(8)	0.10(4)	≤ 0.42
+0.390 ^a	-0.150	-92(2)	12(4)	0.06(2)	0.7(3)
+0.390 ^a	-0.150	-71(3)	18(8)	0.05(2)	1.0(5)
+0.145 ^a	-0.150	-131(1)	8(3)	0.09(3)	≤ 0.71
+0.145 ^a	-0.150	-107(5)	64(9)	0.10(3)	≤ 0.66
+0.063	-0.300	-140(2)	38(4)	0.12(3)	≤ 0.40
-0.460	-0.300	-118(1)	22(1)	0.60(4)	0.09(2)
-0.829	-0.300	-123(3)	39(9)	0.10(3)	≤ 0.66
-0.906	-0.150	-133(1)	22(3)	0.19(5)	≤ 0.27
-0.906	-0.150	-169(2)	27(8)	0.11(5)	≤ 0.40
+0.945 ^a	-0.050	-81(2)	14(3)	0.23(8)	0.3(2)
+0.945 ^a	-0.050	-65(2)	14(3)	0.22(8)	0.18(9)
N					
-1.102	-0.250	-44(2)	39(4)	0.21(5)	0.26(11)
-1.189	-0.300	-4(1)	13(4)	0.19(4)	≤ 0.58
-1.189	-0.300	-42(1)	41(3)	0.43(4)	0.21(5)
-1.189	-0.300	-103(3)	15(6)	0.08(4)	≤ 1.08
-1.211	-0.150	-118(2)	25(2)	0.50(4)	0.26(6)
-1.211	-0.150	-42(1)	34(2)	0.46(4)	0.28(6)